

# A Reduced Surface-Wave Twin Arc-Slot Antenna for Millimeter-Wave Applications

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**Abstract**—A simple, uniplanar, reduced surface-wave twin arc-slot antenna element for millimeter-wave (mm-wave) integrated circuit antenna systems is presented. This design allows for almost complete surface-wave cancellation and resonance to be achieved simultaneously. Measured far-field radiation patterns on an electrically thick substrate are in very good agreement with theory and show little evidence of surface-wave propagation.

**Index Terms**—Coplanar radiation efficiency, integrated antennas, slot antennas, surface waves.

## I. INTRODUCTION

PRINTED slot antennas are useful at millimeter-wave frequencies because of their simplicity, low profile, low cost, light weight, and their ease of integration with electronics. A single-slot antenna element printed on a dielectric substrate at millimeter-wave frequencies suffers greatly from surface-wave losses because typical substrates become electrically thick at these frequencies. The surface-wave losses can be reduced by using broadside linear twin-slot elements to achieve phase cancellation of the dominant  $TM_0$  mode, but only in the broadside direction [1].

A novel twin arc-slot antenna element is proposed here that is shown to reduce the surface-wave fields over a wide angular range in the substrate. This work bears some similarity to the reduced surface-wave circular microstrip patch antennas proposed by Jackson *et al.* in [2]. However, a microstrip patch that excites no surface-waves is larger than a resonant patch, and complex inhomogeneous substrates were proposed in [2] in order to allow complete surface-wave cancellation and the required resonance to be achieved simultaneously. With the proposed twin arc-slot element, achieving nearly complete surface-wave cancellation and resonance simultaneously is made simple by allowing the radius and length of the arc-slot apertures to be chosen independently.

## II. THEORY

The twin arc-slot element geometry is shown in Fig. 1. The arc-slots are printed on a substrate having relative permittivity  $\epsilon_r$  and thickness  $h$ . To avoid excitation of higher-order surface-wave modes and to maximize the front-to-back radiated power ratio, the substrate thickness is made equal to a quarter dielectric wavelength (i.e.,  $h = \lambda_d/4$ ) [1]. As a result, nomi-

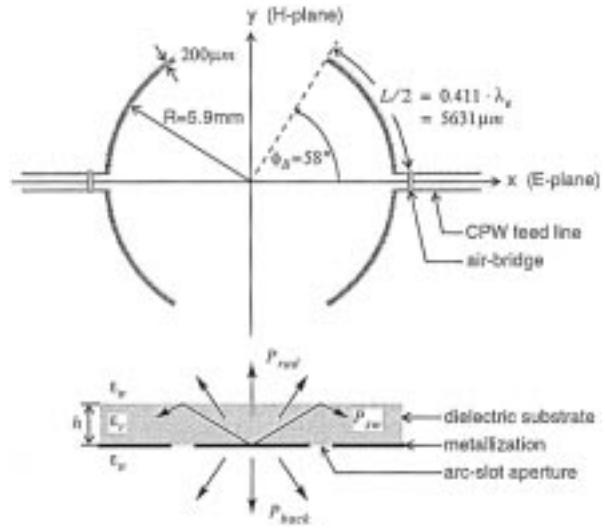


Fig. 1. Geometry of the twin arc-slot antenna.

nally only the dominant  $TM_0$  surface-wave mode is allowed to propagate. The radius of the arc-slots is  $R$ . Each arc-slot is fed at its midpoint by a coplanar waveguide (CPW) feed-line, dividing each arc-slot into two branches of length  $L/2$ , corresponding to a branch arc angle of  $\phi_B$ , as shown in Fig. 1.

The radiation efficiency of a slot antenna is defined by

$$\eta_{\text{rad}} = \frac{P_{\text{rad}}}{P_{\text{rad}} + P_{\text{sw}} + P_{\text{back}}} \quad (1)$$

where  $P_{\text{rad}}$  is the useful power radiated into the air through the substrate,  $P_{\text{sw}}$  is the power coupled into the  $TM_0$  mode (trapped in the substrate), and  $P_{\text{back}}$  is the power radiated directly into the air at the back side of the antenna.

The theoretical analysis of the radiation efficiency of a slot antenna having an arbitrary aperture current distribution has been presented in [3] and [4]. For the twin arc-slot antenna design shown in Fig. 1, the slot aperture is assumed to be very narrow and the magnetic current distribution is assumed to be a sinusoidal one given by

$$\vec{M}(\phi') = \begin{cases} \hat{\phi} \sin[k_e R(\phi_B - |\phi'|)] & 0 < |\phi'| < \phi_B, \\ -\hat{\phi} \sin[k_e R(\phi_B - |\phi' - \pi|)] & 0 < |\phi' - \pi| < \phi_B \end{cases} \quad (2)$$

where  $k_e = 2\pi/\lambda_e$ ,  $\lambda_e = \lambda_0/\sqrt{\epsilon_{\text{eff}}}$  and  $\epsilon_{\text{eff}}$  is the effective dielectric constant, which is approximately equal to  $(\epsilon_r + 1)/2$  for a thick substrate.

The theoretical radiated space-wave power ( $P_{\text{rad}}$  and  $P_{\text{back}}$ ) of the twin arc-slot element can be calculated according to the

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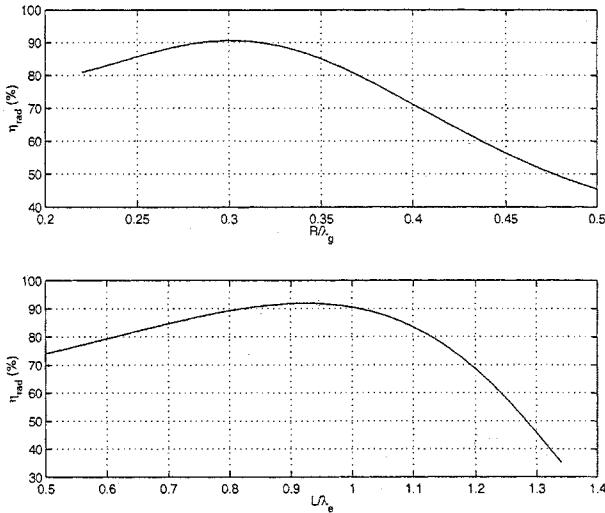


Fig. 2. Radiation efficiency versus  $R$ , at second resonance (i.e.,  $L = \lambda_e$ ), and radiation efficiency versus  $L$ , with  $R = 0.293\lambda_g$ .

method shown in [3] and [4]. The power coupled to the  $\text{TM}_0$  surface-wave mode can be expressed in terms of cylindrical  $\text{TM}_{0n}$  waves emanating from the center of the arc-slot element and the result is given according to [3] as

$$P_{\text{sw}} = P_{\text{TM}_0} = \frac{\omega \epsilon_0 \epsilon_r}{2h_{\text{eff}}} \sum_{n=1}^{+\infty} |C_n|^2 \quad (3)$$

where  $h_{\text{eff}}$  is the effective height of the grounded dielectric substrate corresponding to TM modes, and  $C_n$  is the coupling coefficient of the  $\text{TM}_{0n}$  mode given by

$$C_n = \frac{2[(-1)^n - 1]k_e R^2}{n^2 - (k_e R)^2} \times [\cos(k_e R \phi_B) - \cos(n \phi_B)] [J'_1(\beta_{\text{TM}_0} R)] \quad (4)$$

where  $\beta_{\text{TM}_0} = 2\pi/\lambda_g$  is the propagation constant of the  $\text{TM}_0$  mode, and  $\lambda_g$  is the guide wavelength of the  $\text{TM}_0$  mode.

Only the first few low-order terms in the summation of (3) are significant, and (4) vanishes for  $n = \text{even}$ . If  $R$  is chosen such that  $J'_1(\beta_{\text{TM}_0} R) = 0$  (i.e.,  $R = x'_1/\beta_{\text{TM}_0} = 0.293\lambda_g$ ,  $x'_1 = 1.8412$ , where  $x'_1 = 1.8412$  is the first root of  $J'_1(x)$ ), then the most significant term  $n = 1$  in (3) vanishes, and the only significant remaining terms are those corresponding to  $n = 3$  and  $n = 5$ . These terms can be minimized by properly choosing the angle  $\phi_B$  (i.e., by properly choosing the total arc-slot length  $L$ ). As a result, the surface-wave power is minimized and the corresponding radiation efficiency can be maximized.

### III. NUMERICAL AND EXPERIMENTAL RESULTS

To validate the theory, a substrate having  $\epsilon_r = 10.2$  and  $h = \lambda_d/4$  is considered. The theoretical radiation efficiency is plotted versus  $R$  at second resonance (i.e.,  $L = \lambda_e$ ) and is shown in Fig. 2. A maximum radiation efficiency of 90.7% is achieved near  $R = 0.293\lambda_g$ . Also plotted in Fig. 2 is the radiation efficiency versus  $L$  with  $R = 0.293\lambda_g$ . The efficiency may be increased slightly to 91.9% by reducing  $L$  to approximately 92% of the second resonant length. The efficiency variation is

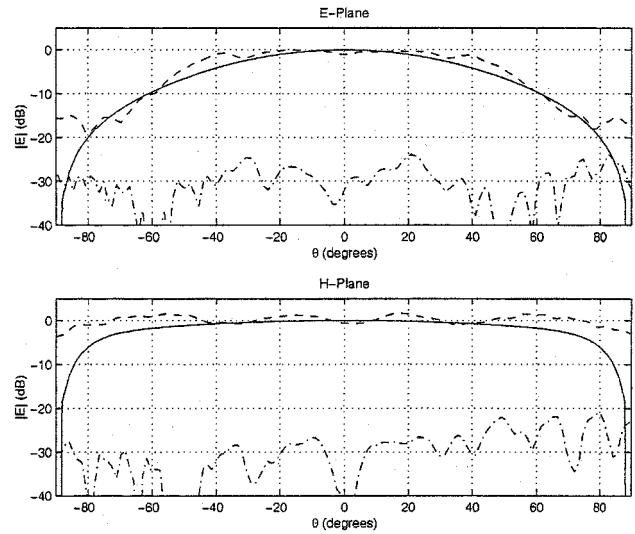


Fig. 3. Far-field  $E$ - and  $H$ -plane radiation patterns at 8.6 GHz.—: Theoretical co-polar; ——: Experimental co-polar; and - - -: Experimental cross-polar.

small for  $L$  between approximately 80% and 100% of the resonant length, and, in practice, resonance typically occurs for a slot length that is slightly shorter than that predicted by transmission line theory [5]. Note that 99.96% of the total power is radiated into the air, but 8.77% of this is lost to  $P_{\text{back}}$ .

The twin arc-slot antenna element shown in Fig. 1 was designed using HP-ADS assuming a *Rogers RT/duroid 6010LM* substrate having  $\epsilon_r = 10.2$  and  $h = 2.54 \text{ mm} = \lambda_d/4$ . The actual size of the substrate was  $12.5 \text{ cm} \times 12.5 \text{ cm}$ . The antenna was fabricated using a wet chemical etching technique. A  $180^\circ$  hybrid coupler was used to provide the required  $180^\circ$  signal phase difference between the two CPW feed-lines so that the arc-slot aperture fields are excited in phase.

The best measured far-field  $E$ - and  $H$ -plane radiation patterns were obtained at a frequency of 8.6 GHz, and are shown in Fig. 3 with the corresponding theoretical patterns predicted using the method-of-moments. All patterns are normalized such that the co-polar components are 0 dB at broadside. The measured co-polar patterns are in very good agreement with theory and show no appreciable rippling effects indicative of surface-wave diffraction from the edges of the finite substrate. Some rippling effects observed may be attributable to the residual excitation of the  $\text{TM}_0$  mode due to fabrication tolerances, lateral-waves which propagate with the free-space wavenumber  $k_0$ , and the fact that the substrate thickness is close to the onset of exciting the  $\text{TE}_0$  mode. The  $\text{TE}_0$  mode close to cutoff can contribute a  $1/r$ -varying space-wave propagating along the substrate. Theoretically, there is no cross polarization in the two principal planes, and the measured cross polarization levels are sufficiently low, being below  $-25 \text{ dB}$  at most angles (see Fig. 3). Although not shown here, the measured patterns in the  $\phi = 45^\circ$  plane also matched well with theory.

### IV. CONCLUSION

A novel reduced surface-wave twin arc-slot antenna element has been proposed, designed, fabricated, and tested. The radius and length of the twin arc-slots may be chosen independently to

achieve effective surface-wave cancellation and resonance simultaneously. The theoretical radiation efficiency of the twin arc-slot antenna element is significantly greater than that of any other planar slot antenna design in the literature. Indeed, the proposed twin arc-slot allows more than 99% of the total power radiated into the air on electrically thick substrates. The measured far-field radiation patterns are in very good agreement with theory, and show no significant ill effects attributable to surface-wave propagation.

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